

Unit - I

2019 Equation of pressure:-

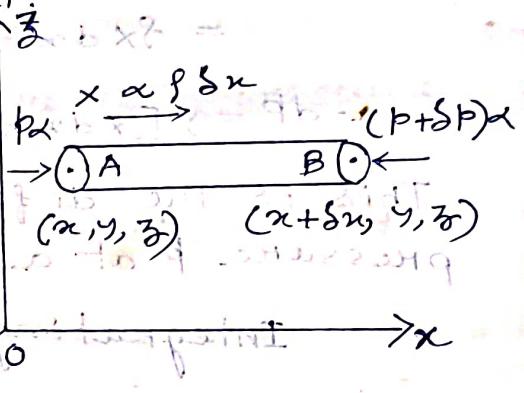
✓ A mass of fluid is at rest under the action of given forces; to obtain the eqn which determines the pressure at any pt. of the fluid.

(6, 00, 96, 09) M

As question:- A mass of fluid is at rest under the action of given forces. obtain in rectangular cartesian co-ordinates, the eqn which determines the pressure at any point.

Let  $A(x, y, z)$  be a pt. in the fluid, referred to rectangular axes.

Let  $B(x + \delta x, y, z)$  be another point very near  $A$ , such that  $AB$  is parallel to  $x$ -axis.



Let us construct a small cylinder about  $AB$  as axis with its plane ends  $\perp$  to  $AB$ .

Let  $\alpha$  be the area of the plane ends of the cylinder.

Also, let  $p$  and  $p + \delta p$  be the pressure at  $A$  and  $B$ , respectively.

Therefore thrust at the plane end at  $A$  is  $p\alpha$  and that at  $B$  is  $(p + \delta p)\alpha$ .  
Let  $\rho$  be the mean density of the liquid in the cylinder  $AB$ , then mass of the cylinder =  $\rho\alpha\delta x$ .

Now, let  $x, y, z$  be the components of given forces along  $x$ ,  $y$ ,  $z$ -axes respectively per unit mass, then  $x\alpha g\delta x$  is the force on the cylinder  $\parallel$  to the  $x$ -axis.

Thus the cylinder is in equilibrium under the forces

- i)  $p\alpha$  along  $AB$
- ii)  $x\alpha g\delta x$  along  $AB$ , and
- iii)  $(p + \delta p)\alpha$  along  $BA$ .

$\therefore$  For the equilibrium of the cylinder, we have

$$(p + \delta p)x = px + x\delta p$$

$$\text{or } \delta p = gx \delta x \text{ or } \frac{\delta p}{\delta x} = gx$$

Taking limit as  $\delta x$  tends to zero, we get

$$\frac{\partial p}{\partial x} = gx$$

By, we have  $\frac{\partial p}{\partial y} = gy$  and  $\frac{\partial p}{\partial z} = gz$

Now,  $dP = \frac{\partial P}{\partial x} dx + \frac{\partial P}{\partial y} dy + \frac{\partial P}{\partial z} dz$  [as  $P$  is a function of  $x, y, z$ ]

$$= gx dx + gy dy + gz dz$$

This is the differential eqn determining pressure  $p$  at a point  $(x, y, z)$ .

## \*<sup>2015</sup> Necessary and sufficient conditions for equilibrium

of a fluid:-

Ex) State and establish the NASC that must be satisfied by a given system of forces  $x, y, z$  acting per unit mass of a fluid and  $\Pi$ , to the axes of a rectangular cartesian co-ordinate system so that the fluid may maintain equilibrium.

Statement:- The NASC for a fluid may maintain equilibrium, which must be satisfied under a given system of forces  $x, y, z$  acting per unit mass of the fluid and  $\Pi$ , to the rectangular

Co-ordinate axes is

$$x \left( \frac{\partial \gamma}{\partial z} - \frac{\partial z}{\partial y} \right) + \gamma \left( \frac{\partial z}{\partial x} - \frac{\partial x}{\partial z} \right) + z \left( \frac{\partial x}{\partial y} - \frac{\partial y}{\partial x} \right) = 0 \quad (1)$$

Proof:- Necessary part :-

Let us take  $x$ -axis and  $y$ -axis be horizontal and  $z$ -axis vertically upward.

Let the fluid be in equilibrium.

Let  $P(x, y, z)$  be a ph. in the fluid.

Let  $\varphi(x + \delta x, y, z)$  be another ph. very near to  $P$  such that  $P \parallel x$ -axis.

Construct a small cylinder about  $P\varphi$  as axis with its plane end  $S$   $\perp$  to  $P\varphi$ ,

Let  $\alpha$  be the area of the plane ends of the cylinder. Also let  $p$  and  $(p + \delta p)$  be the pressure at  $P$  and  $\varphi$  respectively.

There force thrust at the plane end at  $P$  is  $p\alpha$ , along  $P\varphi$

and that at  $\varphi$  is  $(p + \delta p)\alpha$ , along  $\varphi P$ .

Now, let  $\rho$  be the mean density of the liquid in the cylinder  $P\varphi$ , then mass of the cylinder  $= \rho\alpha\delta x$ .

Let  $x, y, z$  be the components of given forces

per unit mass along  $x, y, z$ -axes respectively, then the force on the cylinder  $\parallel$  to  $x$ -axis is  $x\alpha\delta x$ , along  $P\varphi$ . For equilibrium,

$$(p + \delta p)\alpha = p\alpha + x\alpha\delta x$$

$$\Rightarrow \delta p = x\delta x$$

$$\Rightarrow \frac{\delta p}{\delta x} = x$$

Taking limit as  $\delta x \rightarrow 0$ , we get

$$\frac{dp}{dx} = x$$

if, we have

$$\frac{\partial p}{\partial y} = \rho T_n \rightarrow (A)$$

$$\text{and } \frac{\partial p}{\partial z} = \rho z$$

we assume that  $p$  is continuously differentiable.

$$\therefore \frac{\partial^2 p}{\partial x \partial y} = \frac{\partial^2 p}{\partial y \partial x}, \quad \frac{\partial^2 p}{\partial x^2} = \frac{\partial^2 p}{\partial y^2}, \quad \frac{\partial^2 p}{\partial x \partial z} = \frac{\partial^2 p}{\partial y \partial z} \text{ and similarly}$$

$$\text{Now, } \frac{\partial^2 p}{\partial z \partial y} = \frac{\partial^2 p}{\partial y \partial z}$$

$$\Rightarrow \frac{\partial}{\partial y} \left( \frac{\partial p}{\partial z} \right) = \frac{\partial}{\partial z} \left( \frac{\partial p}{\partial y} \right)$$

$$\Rightarrow \frac{\partial}{\partial y} (\delta z) = \frac{\partial}{\partial z} (\delta y)$$

$$\Rightarrow \gamma \frac{\partial z}{\partial y} + \delta \frac{\partial z}{\partial y} = \gamma \frac{\partial y}{\partial z} + \delta \frac{\partial y}{\partial z}$$

$$\Rightarrow \delta \left( \frac{\partial z}{\partial y} - \frac{\partial y}{\partial z} \right) = \gamma \frac{\partial y}{\partial z} - \gamma \frac{\partial z}{\partial y} \rightarrow \textcircled{1}$$

likewise;  $\frac{\partial^2 p}{\partial z \partial x} = \frac{\partial^2 p}{\partial x \partial z} \Rightarrow \delta \left( \frac{\partial z}{\partial x} - \frac{\partial x}{\partial z} \right) = \gamma \frac{\partial x}{\partial z} - \gamma \frac{\partial z}{\partial x} \rightarrow \textcircled{2}$

$$\frac{\partial^2 p}{\partial x \partial y} = \frac{\partial^2 p}{\partial y \partial x} \Rightarrow \delta \left( \frac{\partial x}{\partial y} - \frac{\partial y}{\partial x} \right) = \gamma \frac{\partial y}{\partial x} - \gamma \frac{\partial x}{\partial y} \rightarrow \textcircled{3}$$

Now, multiplying  $\textcircled{1}, \textcircled{2}, \textcircled{3}$  by  $x, y, z$  and adding, we get

$$\times \left( \frac{\partial y}{\partial z} - \frac{\partial z}{\partial y} \right) + \gamma \left( \frac{\partial z}{\partial x} - \frac{\partial x}{\partial z} \right) + \gamma \left( \frac{\partial x}{\partial y} - \frac{\partial y}{\partial x} \right) = 0$$

Hence the condition is necessary.

Sufficient condition:-

$$\text{Let } \times \left( \frac{\partial y}{\partial z} - \frac{\partial z}{\partial y} \right) + \gamma \left( \frac{\partial z}{\partial x} - \frac{\partial x}{\partial z} \right) + \gamma \left( \frac{\partial x}{\partial y} - \frac{\partial y}{\partial x} \right) = 0 \rightarrow \textcircled{4}$$

Now,

$$\frac{\partial}{\partial z} (\delta y) - \frac{\partial}{\partial y} (\delta z) = \delta \frac{\partial y}{\partial z} + \gamma \frac{\partial z}{\partial z} - \gamma \frac{\partial z}{\partial y} - \gamma \frac{\partial y}{\partial z} \rightarrow \textcircled{5}$$

$$\frac{\partial}{\partial x} (\delta z) - \frac{\partial}{\partial z} (\delta x) = \delta \frac{\partial z}{\partial x} + \gamma \frac{\partial x}{\partial z} - \gamma \frac{\partial x}{\partial z} - \gamma \frac{\partial z}{\partial x} \rightarrow \textcircled{6}$$

$$\frac{\partial}{\partial y} (\delta x) - \frac{\partial}{\partial x} (\delta y) = \delta \frac{\partial x}{\partial y} + \gamma \frac{\partial y}{\partial x} - \gamma \frac{\partial y}{\partial x} - \gamma \frac{\partial x}{\partial y} \rightarrow \textcircled{7}$$

Now, multiplying  $\textcircled{5}, \textcircled{6}, \textcircled{7}$  by  $\delta x, \delta y, \delta z$ , respectively and adding, we get

$$\begin{aligned} & \delta x \left[ \frac{\partial}{\partial z} (\delta y) - \frac{\partial}{\partial y} (\delta z) \right] + \delta y \left[ \frac{\partial}{\partial x} (\delta z) - \frac{\partial}{\partial z} (\delta x) \right] \\ & + \delta z \left[ \frac{\partial}{\partial y} (\delta x) - \frac{\partial}{\partial x} (\delta y) \right] \\ &= \delta^2 \left[ \times \left( \frac{\partial y}{\partial z} - \frac{\partial z}{\partial y} \right) + \gamma \left( \frac{\partial z}{\partial x} - \frac{\partial x}{\partial z} \right) + \gamma \left( \frac{\partial x}{\partial y} - \frac{\partial y}{\partial x} \right) \right] \\ & + \delta \left[ \times \left( \gamma \frac{\partial y}{\partial z} - \gamma \frac{\partial z}{\partial y} \right) + \gamma \left( \gamma \frac{\partial z}{\partial x} - \gamma \frac{\partial x}{\partial z} \right) \right. \\ & \quad \left. + \gamma \left( \gamma \frac{\partial x}{\partial y} - \gamma \frac{\partial y}{\partial x} \right) \right] \\ &= \delta^2 (0 + 0 + 0) = 0 \end{aligned}$$

Thus, ~~the condition~~  $(\rho_x)dx + (\rho_Y)dy + (\rho_Z)dz$  is a perfect differential of some function.

$$\therefore \rho_x dx + \rho_Y dy + \rho_Z dz = dp \text{ (say)}$$

$$\Rightarrow \rho(xdx + Ydy + Zdz) = dp - \frac{\partial p}{\partial x} dx + \frac{\partial p}{\partial Y} dy + \frac{\partial p}{\partial Z} dz$$

$$\Rightarrow \frac{\partial p}{\partial x} = \rho_x, \frac{\partial p}{\partial Y} = \rho_Y, \frac{\partial p}{\partial Z} = \rho_Z$$

But these are clearly the eqns of equilibrium of a fluid. Thus if ④ holds, the eqns of equilibrium follow.

Hence the condition is sufficient.

2018 (89, 94) M

Ex Deduce that for a homogeneous fluid at rest under gravity, the free surface is horizontal.

Soln: We have the pressure eqn is

$$[(x) \frac{\partial p}{\partial x} + (y) \frac{\partial p}{\partial y} + (z) \frac{\partial p}{\partial z}] = \rho g$$

where  $x, y$  are horizontal forces and  $z$  is the vertically upward force. Here  $x = 0 = y$  and  $z = -g$ . Hence (i) becomes

$$\frac{\partial p}{\partial z} = -\rho g$$

Integrating, we get  $p = K - \frac{1}{2} \rho g z^2$

$p = K - \rho g z$  ( $\because \rho$  is const. because the fluid is homogeneous)  $\rightarrow$  (ii)

Let  $p = p_0$  (const) at the origin, (i.e.  $z = 0$ )

$$\therefore P_0 = k$$

$$P = P_0 - \rho g z \rightarrow \text{(iii)}$$

At the free surface,  $P = 0$

eq<sup>n</sup> (iii) becomes,  $P_0 = \rho g z$

$$\Rightarrow z = \frac{P_0}{\rho g} = h \text{ (say)}$$

which is the horizontal plane

Homogeneous liquids:- We know that the pressure at a pt. is given by,

$$dp = \rho(x dx + y dy + z dz) \rightarrow \textcircled{1}$$

where  $\rho$  is the density of the ~~fluid~~ liquid at the pt. and  $x, y, z$  are the component forces per unit mass.

Now, if the liquid is homogeneous, then it's constraint, thus from  $\textcircled{1}$ , we get

$$x dx + y dy + z dz = \frac{1}{\rho} dp = d\left(\frac{p}{\rho}\right) \rightarrow \textcircled{2}$$

i.e  $x dx + y dy + z dz$  is a perfect differential. But if  $x dx + y dy + z dz$  is a perfect differential, then the system of forces is said to be conservative. Therefore a homogeneous liquid will be in equilibrium only when the system of forces is conservative. Thus in this case, we may write,

$$x dx + y dy + z dz = -dv, \quad v \text{ being potential function.}$$

$$\text{Or } \frac{dp}{\rho} = -dv \text{ from } \textcircled{2}$$

Integrating, we get

$$\frac{p}{\rho} + v = c, \quad c \text{ being constant of integration.}$$

Heterogeneous fluids:-

In a heterogeneous fluid,  $\rho$  varies, i.e.  $\rho$  is a function of independent variables  $x, y, z$ . In this case the system of forces  $x, y, z$  may maintain equilibrium if

$$\frac{\partial}{\partial y} (\rho z) = \frac{\partial}{\partial z} (\rho y), \quad \frac{\partial}{\partial z} (\rho x) = \frac{\partial}{\partial x} (\rho z),$$

$$\frac{\partial}{\partial x} (\rho y) = \frac{\partial}{\partial y} (\rho x), \quad \text{from } \textcircled{1} \text{ in case.}$$

## Surfaces of equal pressure:-

For a liquid at rest in equilibrium, the pressure at any pt. is given by,

$$dp = \rho(xdx + ydy + zdz)$$

Integrating, we get

$$p = \phi(x, y, z) \text{ (say)}$$

If  $p = \text{constant} = c$  (say), then we have a surface.

$$\phi(x, y, z) = c \rightarrow \textcircled{O}$$

at every pt. of which the pressure is constant and equal to  $c$ .

such a surface is called surface of equal pressure.

For different values of  $c$ , we obtain different surfaces, at every pt. of these surfaces the pressure is same.

Thus as  $c$  takes different values,  $\textcircled{O}$  represents surfaces of equal pressure.

The external surface, one surface is obtained by making  $p$  equal to the pressure extend to the fluid. If external pressure is zero, the free surface is given by

$$\phi(x, y, z) = 0$$

Line of force:- A line of force is a curve (i.e. line) such that the tangent at every pt. of

which is in the direction of the resultant force at that point.

As question :- Define a surface of equal

pressure and a line of force for a fluid

in equilibrium under a given system of forces.

## Curves of equal pressure and density :-

2016

If a fluid is at rest under the forces  $x, y, z$  per unit mass, then to find the differential eqns of the curves of equal pressure and equal density.

(07,02) M

As question:-

Obtain the differential eqns of the curves of equal pressure and density in the case of other fluid which is acted on by forces  $x, y, z$  per unit mass along the rectangular axes.

Soln:- Let the fluid be heterogeneous (so that the density is a function of  $x, y, z$ ) and incompressible. Then the surfaces of equal pressure are given by,  $p = \text{constant}$ . i.e.  $dp = 0$

$$\text{or } \varphi(x dx + y dy + z dz) = 0 \rightarrow \textcircled{1}$$

And surfaces of equal density are given by,  $\varphi = \text{constant}$ .

$$\text{or } \frac{\partial \varphi}{\partial x} dx + \frac{\partial \varphi}{\partial y} dy + \frac{\partial \varphi}{\partial z} dz = 0 \rightarrow \textcircled{2}$$

The curves of intersection of  $\textcircled{1}$  and  $\textcircled{2}$  are the curves of equal pressure and density.

Solving  $\textcircled{1}$  and  $\textcircled{2}$ , we get

$$\frac{dx}{Y \frac{\partial p}{\partial z} - Z \frac{\partial p}{\partial y}} = \frac{dy}{Z \frac{\partial p}{\partial x} - X \frac{\partial p}{\partial z}} = \frac{dz}{X \frac{\partial p}{\partial y} - Y \frac{\partial p}{\partial x}} \rightarrow \textcircled{3}$$

We assume that  $p$  is continuously differentiable

$$\therefore \frac{\partial^2 p}{\partial y \partial z} = \frac{\partial^2 p}{\partial z \partial y}, \frac{\partial^2 p}{\partial z \partial x} = \frac{\partial^2 p}{\partial x \partial z}, \frac{\partial^2 p}{\partial x \partial y} = \frac{\partial^2 p}{\partial y \partial x}$$

$$\text{Now, } \frac{\partial^2 p}{\partial y \partial z} - \frac{\partial^2 p}{\partial z \partial y} \Rightarrow \frac{\partial}{\partial y} \left( \frac{\partial p}{\partial z} \right) = \frac{\partial}{\partial z} \left( \frac{\partial p}{\partial y} \right)$$

$$\Rightarrow \frac{\partial}{\partial y} (yz) = \frac{\partial}{\partial z} (zy)$$

$$\Rightarrow z \frac{\partial z}{\partial y} + y \frac{\partial z}{\partial y} = Y \frac{\partial z}{\partial z} + Z \frac{\partial z}{\partial y} \Rightarrow -y = -z$$

$$\Rightarrow Y \frac{\partial z}{\partial z} - Z \frac{\partial z}{\partial y} = -y \left[ \frac{\partial z}{\partial z} - \frac{\partial z}{\partial y} \right] \rightarrow \textcircled{4}$$

$$i) \quad \frac{\partial p}{\partial x} - \gamma \frac{\partial p}{\partial z} = -p \left[ \frac{\partial z}{\partial x} - \frac{\partial x}{\partial z} \right] \rightarrow ⑤$$

$$\text{and } \gamma \frac{\partial p}{\partial y} - \gamma \frac{\partial p}{\partial x} = -p \left[ \frac{\partial x}{\partial y} - \frac{\partial y}{\partial x} \right] \rightarrow ⑥$$

Putting, ④, ⑤, ⑥ in ③, we get

$$\frac{dx}{-p \left[ \frac{\partial y}{\partial z} - \frac{\partial z}{\partial y} \right]} = \frac{dy}{-p \left[ \frac{\partial z}{\partial x} - \frac{\partial x}{\partial z} \right]} = \frac{dz}{-p \left[ \frac{\partial x}{\partial y} - \frac{\partial y}{\partial x} \right]}$$

$$\Rightarrow \frac{dx}{\frac{\partial y}{\partial z} - \frac{\partial z}{\partial y}} = \frac{dy}{\frac{\partial z}{\partial x} - \frac{\partial x}{\partial z}} = \frac{dz}{\frac{\partial x}{\partial y} - \frac{\partial y}{\partial x}}$$

which are the differential eqns of. the curves  
of equal pressure and density.

(Q6) ~~2017 Ex)~~ A liquid of given volume  $V$  is at rest under the forces,  $\vec{F} = -\frac{u_x}{a^2} \hat{x} + -\frac{u_y}{b^2} \hat{y} + -\frac{u_z}{c^2} \hat{z}$

Find the pressure at any pt.  $x$  of the liquid and  
the surfaces of equal pressure.

Soln:- The pressure eqn is,

$$dp = \rho (x dx + y dy + z dz) \rightarrow (i)$$

$$= \rho \left( -u \frac{x}{a^2} dx - u \frac{y}{b^2} dy - u \frac{z}{c^2} dz \right)$$

$$= -\rho u \left( \frac{x}{a^2} dx + \frac{y}{b^2} dy + \frac{z}{c^2} dz \right)$$

Integrating;

$$p = -\rho u \left[ \frac{x^2}{2a^2} + \frac{y^2}{2b^2} + \frac{z^2}{2c^2} \right] + c_1, c_1 = \text{arbitrary constant}$$

$$\Rightarrow p = -\frac{\rho u}{2} \left[ \frac{x^2}{a^2} + \frac{y^2}{b^2} + \frac{z^2}{c^2} \right] + c_1 \rightarrow (ii)$$

Now, on the free surface  $P=0$ ,

$$\Rightarrow 0 = -\frac{g}{2} \mu \left[ \frac{x^2}{a^2} + \frac{y^2}{b^2} + \frac{z^2}{c^2} \right] + C_1$$

$$\Rightarrow C_1 = \frac{1}{2} g \mu \left[ \left( \frac{x}{a} \right)^2 + \left( \frac{y}{b} \right)^2 + \left( \frac{z}{c} \right)^2 \right]$$

$$\Rightarrow \left( \frac{x}{a} \right)^2 + \left( \frac{y}{b} \right)^2 + \left( \frac{z}{c} \right)^2 = \frac{2C_1}{g\mu} = \gamma^2 \text{ (say)} \rightarrow (iii)$$

$$\Rightarrow \left( \frac{x}{a\gamma} \right)^2 + \left( \frac{y}{b\gamma} \right)^2 + \left( \frac{z}{c\gamma} \right)^2 = 1$$

which clearly represents an ellipsoid. But we know that the volume  $V$  of the ellipsoid is given by;

$$V = \frac{4}{3} \pi (a\gamma)(b\gamma)(c\gamma) \rightarrow \text{by the question, } V \text{ is constant (given)}$$

$$\Rightarrow \frac{3V}{4\pi} = abc\gamma^3$$

$$\Rightarrow \gamma^3 = \frac{3V}{4\pi abc \mu}$$

$$\Rightarrow \gamma = \left[ \frac{3V}{4\pi abc} \right]^{1/3}$$

$$\therefore (iii) \Rightarrow \frac{2C_1}{g\mu} = \gamma^2 = \left[ \frac{3}{4\pi} \cdot \frac{V}{abc} \right]^{2/3}$$

$$\Rightarrow C_1 = \frac{1}{2} g \mu \left[ \frac{3V}{4\pi abc} \right]^{2/3}$$

$$\therefore (ii) \Rightarrow P = -\frac{1}{2} g \mu \left[ \frac{x^2}{a^2} + \frac{y^2}{b^2} + \frac{z^2}{c^2} \right] + \frac{1}{2} g \mu \left[ \frac{3V}{4\pi abc} \right]^{2/3}$$

$$= \frac{1}{2} g \mu \left[ \frac{3V}{4\pi abc} \right]^{2/3} - \frac{1}{2} g \mu \left[ \frac{x^2}{a^2} + \frac{y^2}{b^2} + \frac{z^2}{c^2} \right]$$

This gives pressure at any point.

Again, the equiprессure surface is given by,

$$\Rightarrow g \left[ x dx + y dy + z dz \right] = 0$$

$$\Rightarrow \frac{x}{a^2} dx + \frac{y}{b^2} dy + \frac{z}{c^2} dz = 0$$

Integrating  $\Rightarrow \frac{x^2}{a^2} + \frac{y^2}{b^2} + \frac{z^2}{c^2} = k$ , where  $k$  is an

arbitrary constant which represents a family of surfaces of equal pressure.

<sup>(09) M</sup>  
~~\* Ex~~ Prove that if the forces per unit mass acts parallel to the axes are  $y(a-z)$ ,  $x(a-z)$ ,  $xy$  the surfaces of equal pressure are hyperbolic paraboloids and the curves of equal pressure and density are rectangular hyperboloids.

Soln:- As given, we have

$$x = y(a-z), Y = z(a-x), \frac{z}{x} = xy$$

so that,  $\frac{\partial x}{\partial y} = a-z, \frac{\partial x}{\partial z} = -y, \frac{\partial y}{\partial x} = a-z, \frac{\partial y}{\partial z} = -x$

$$\frac{\partial Y}{\partial x} = -z, \frac{\partial Y}{\partial z} = y \text{ and } \frac{\partial z}{\partial y} = x$$

The condition of equilibrium is,

$$x\left(\frac{\partial Y}{\partial z} + \frac{\partial z}{\partial y}\right) + Y\left(\frac{\partial z}{\partial x} + \frac{\partial x}{\partial z}\right) + z\left(\frac{\partial x}{\partial y} + \frac{\partial y}{\partial x}\right) = 0$$

$$\text{i.e. } y(a-z)(-x-z) + z(a-x)(y+y) + xy[a-z-(a-z)]$$

which is clearly satisfied.

Now, surfaces of equal pressure are given by  $p = \text{constant}$  i.e.  $dp = 0$  or

$$xdx + Ydy + zdz = 0$$

$$\text{or } y(a-z)dx + z(a-x)dy + xydz = 0 \quad (\text{iii})$$

Putting values of  $x, Y, z$ ,

Dividing by  $xy(a-z)$ , the surfaces of..

equal pressure are given by,

$$\frac{dx}{x} + \frac{dy}{y} + \frac{dz}{a-z} = 0$$

Integrating  $\log x + \log y + \log(a-z) = \log c$  or

$\frac{xy}{a-z} = c$  which are clearly hyperbolic paraboloids.

Again the curves of equal pressure and equal density are given by,

$$\frac{dx}{\frac{\partial Y}{\partial z} - \frac{\partial z}{\partial y}} = \frac{dy}{\frac{\partial z}{\partial x} - \frac{\partial x}{\partial z}} = \frac{dz}{\frac{\partial x}{\partial y} - \frac{\partial y}{\partial x}}$$

$$\text{or } \frac{dx}{-z} = \frac{dy}{y} = \frac{dz}{0} \quad (\text{iv})$$

The last fraction gives  $dz = 0$  or  $z = \text{constant}$ .  
And the first two fractions are

$\frac{dx}{x} + \frac{dy}{y} = 0$  or  $xy = \text{constant}$ . To understand this  
thus the curves of equal pressure and equal  
density are given by (ii) part -

$xy = \text{constant}$ ,  $z = \text{constant}$ ,  
which are clearly rectangular hyperbolae.

(09) M

Ex). A closed tube in the form of an ellipse with its major axis vertical, is filled with three different liquids of densities  $\rho_1, \rho_2$  and  $\rho_3$  respectively. If the distances of the surfaces of separation from either focus be  $M_1, M_2, M_3$  respectively. Prove that-

$$M_1(\rho_2 - \rho_3) + M_2(\rho_3 - \rho_1) + M_3(\rho_1 - \rho_2) = 0$$

Soln:- Let  $s$  be the focus and  $zz'$  be the directrix of the ellipse. Let  $A_1A_2, A_2A_3, A_3A_1$  be the liquids of densities  $\rho_3, \rho_1, \rho_2$  respectively.

Let  $SA_1 = k_1, SA_2 = k_2, SA_3 = k_3$ .  
Let  $A_1M_1, A_2M_2, A_3M_3$  be the  $\perp^s$  on the directrix  $zz'$ .

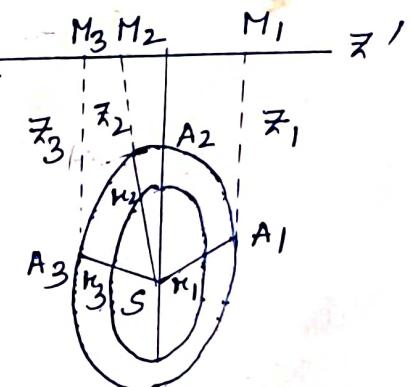
$$\text{Let } A_1M_1 = z_1, A_2M_2 = z_2,$$

$$A_3M_3 = z_3.$$

From the properties of ellipse, we know that the distance of any pt. on the ellipse from the focus = ex the  $\perp^r$  distance of the pt. from the directrix where  $e$  is the eccentricity of the ellipse.

$$\text{Hence } k_1 = e z_1, k_2 = e z_2, k_3 = e z_3$$

$$\Rightarrow z_1 = \frac{k_1}{e} \quad \Rightarrow z_2 = \frac{k_2}{e} \quad \Rightarrow z_3 = \frac{k_3}{e}$$



Let the pressure at the pt. of separation of the liquids  $A_1, A_2, A_3$ , be  $p_1, p_2, p_3$  respectively.

$$p_2 - p_1 = \gamma_3 g (\bar{z}_2 - \bar{z}_1)$$

$$p_3 - p_2 = \gamma_1 g (\bar{z}_3 - \bar{z}_2)$$

$$p_1 - p_3 = \gamma_2 g (\bar{z}_1 - \bar{z}_3)$$

Adding all these and cancelling  $g$ , we get

$$0 = \gamma_3 (\bar{z}_2 - \bar{z}_1) + \gamma_1 (\bar{z}_3 - \bar{z}_2) + \gamma_2 (\bar{z}_1 - \bar{z}_3)$$

$$= \gamma_3 \left( \frac{\kappa_2}{e} - \frac{\kappa_1}{e} \right) + \gamma_1 \left( \frac{\kappa_3}{e} - \frac{\kappa_2}{e} \right) + \gamma_2 \left( \frac{\kappa_1}{e} - \frac{\kappa_3}{e} \right)$$

$$\Rightarrow (\kappa_2 - \kappa_1) \gamma_3 + (\kappa_3 - \kappa_2) \gamma_1 + (\kappa_1 - \kappa_3) \gamma_2 = 0$$

$$\Rightarrow \kappa_2 \gamma_3 - \kappa_1 \gamma_3 + \kappa_3 \gamma_1 - \kappa_2 \gamma_1 + \kappa_1 \gamma_2 - \kappa_3 \gamma_2 = 0$$

$$\Rightarrow \kappa_1 (\gamma_2 - \gamma_3) + \kappa_2 (\gamma_3 - \gamma_1) + \kappa_3 (\gamma_1 - \gamma_2) = 0$$

Proved

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**Example 3.** If  $X = y(y+z)$ ,  $Y = z(z+x)$ ,  $Z = y(y-x)$ , surfaces of equal pressure are the hyperbolic paraboloids  $y(x+z) = c(y+z)$  and the curves of equal pressure and density are given by  $y(x+z) = \text{const.}$ ,  $(y+z) = \text{const.}$

**Sol.** We know that

$$dp = \rho [X dx + Y dy + Z dz].$$

Now for surfaces of equal pressure,  $p = \text{const.}$ , i.e.,  $dp = 0$ .

Thus surfaces of equal pressure are given by  $dp = 0$

$$X dx + Y dy + Z dz = 0$$

or

$$y(y+z) dx + z(z+x) dy + y(y-x) dz = 0$$

or

$$\frac{dx}{z+x} + \frac{z dy}{y(y+z)} + \frac{(y-x) dz}{(y+z)(z+x)} = 0$$

or

$$\frac{dx}{z+x} + \frac{z dy}{y(y+z)} + \frac{(y+z)-(x+z)}{(y+z)(z+x)} dz = 0,$$

or

$$\text{as } y-x = (y+z)-(x+z)$$

$$\frac{dx}{z+x} + \frac{z dy}{y(y+z)} + \frac{dz}{z+x} - \frac{dz}{y+z} = 0$$

or

$$\frac{dx+dz}{x+z} + \frac{z dy - y dz}{y(y+z)} = 0$$

or

$$\frac{dx+dz}{x+z} + \frac{dy(z+y) - y(dy+dz)}{y(y+z)} = 0$$

or

$$\frac{dx+dz}{x+z} + \frac{dy}{y} - \frac{dy+dz}{y+z} = 0$$

Integrating,  $\log(x+z) + \log y - \log(y+z) = \log c$

or

$$y(x+z) = c(y+z),$$

which is the equation giving surfaces of equal pressure.

Again curves of equal pressure and density are given by

$$\frac{\frac{dx}{\partial Y} - \frac{\partial Z}{\partial y}}{\frac{\partial Z}{\partial z} - \frac{\partial X}{\partial y}} = \frac{\frac{dy}{\partial Z} - \frac{\partial X}{\partial z}}{\frac{\partial X}{\partial y} - \frac{\partial Y}{\partial x}} = \frac{\frac{dz}{\partial X} - \frac{\partial Y}{\partial z}}{\frac{\partial Y}{\partial z} - \frac{\partial X}{\partial z}}$$

or

$$\frac{dx}{(2z+x)-(2y-x)} = \frac{dy}{-y-y} = \frac{dz}{(2y+z)-z}$$

or

$$\frac{dx}{x+z-y} = \frac{dy}{-y} = \frac{dz}{y} \quad \dots(1)$$

From the last two fractions,

$$dy + dz = 0 \quad \text{or} \quad y + z = \text{const.}$$

Also using the multipliers  $y, x+z, y$  respectively, we have

$$y \, dx + (x+z) \, dy + y \, dz = 0$$

or

$$y \, dx + x \, dy + z \, dy + y \, dz = 0$$

or

$$xy + yz = \text{const.} \quad \dots(2)$$

(1) and (2) together represent curves of equal pressure and density

**Example 4.** Show that the forces represented by

$$X = \mu (y^2 + yz + z^2), Y = \mu (z^2 + zx + x^2), Z = \mu (x^2 + xy + y^2)$$

will keep a mass of liquid at rest, if the density  $\propto \frac{1}{(\text{dist.})^2}$  from the plane  $x + y + z = 0$ ;

and the curves of equal pressure and density will be circles.

**Sol.** Let  $d$  be the distance of point  $(x, y, z)$  from the plane  $x + y + z = 0$ ; then

$$d = \frac{x+y+z}{\sqrt{3}}$$

If  $\rho$  be the density and  $\rho \propto \frac{1}{(\text{dist.})^2}$ , then

$$\rho \propto \frac{1}{d^2} \quad \text{or} \quad \rho = \frac{\lambda}{(x+y+z)^2}, \quad \dots(1)$$

where  $\lambda$  is a const, the liquid is heterogeneous.

Now the given forces will keep the heterogeneous fluid in equilibrium, if

$$\frac{\partial}{\partial y} (\rho Z) = \frac{\partial}{\partial z} (\rho Y), \quad \frac{\partial}{\partial x} (\rho Z) = \frac{\partial}{\partial z} (\rho X), \quad \frac{\partial}{\partial y} (\rho X) = \frac{\partial}{\partial x} (\rho Y).$$

We have to show that these conditions are satisfied by the given values of  $X, Y, Z$  when  $\rho$  is given by (1).

$$\begin{aligned} \text{Now } \frac{\partial}{\partial y} (\rho Z) &= \frac{\partial}{\partial y} \left[ \frac{\lambda}{(x+y+z)^2} \times \mu (x^2 + xy + y^2) \right] \\ &= \lambda \mu \frac{(x+2y)(x+y+z)^2 - 2(x+y+z)(x^2 + xy + y^2)}{(x+y+z)^4} \\ &= \lambda \mu \frac{[(x+2y)(x+y+z) - 2(x^2 + xy + y^2)]}{(x+y+z)^3} \\ &= \frac{\lambda \mu}{(x+y+z)^3} [-x^2 + 2yz + xz + xy]. \end{aligned} \quad \dots(2)$$

and

$$\begin{aligned} \frac{\partial}{\partial z} (\rho Y) &= \frac{\partial}{\partial z} \left[ \frac{\lambda}{(x+y+z)^2} \mu (z^2 + zx + x^2) \right] \\ &= \lambda \mu \frac{(2z+x)(x+y+z)^2 - 2(x+y+z)(z^2 + zx + x^2)}{(x+y+z)^4} \\ &= \lambda \mu \frac{[(2z+x)(x+y+z) - 2(z^2 + zx + x^2)]}{(x+y+z)^3} \\ &= \frac{\lambda \mu}{(x+y+z)^3} [-x^2 + zx + 2yz + xy]. \end{aligned} \quad \dots(3)$$

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Thus from (2) and (3), we get  $\frac{\partial}{\partial y} (\rho Z) = \frac{\partial}{\partial z} (\rho Y)$ .

Similarly it can be shown that

$$\frac{\partial}{\partial x} (\rho Z) = \frac{\partial}{\partial z} (\rho X) \text{ and } \frac{\partial}{\partial y} (\rho X) = \frac{\partial}{\partial x} (\rho Y).$$

Thus the given forces will keep the mass of liquid at rest if  $\rho$  is given by (1). Again the curves of equal pressure and density are

$$\frac{dx}{\frac{\partial Y}{\partial z} - \frac{\partial Z}{\partial y}} = \frac{dy}{\frac{\partial Z}{\partial x} - \frac{\partial X}{\partial z}} = \frac{dz}{\frac{\partial X}{\partial y} - \frac{\partial Y}{\partial x}}$$

$$\text{or } \frac{dx}{(x+2z)-(x+2y)} = \frac{dy}{(2x+y)-(y+2z)} = \frac{dz}{(2y+z)-(z+2x)}$$

$$\text{or } \frac{dx}{z-y} = \frac{dy}{x-z} = \frac{dz}{y-x}$$

$$= \frac{dx + dy + dz}{0} = \frac{x dx + y dy + z dz}{0}$$

Now  $dx + dy + dz = 0$  gives  $x + y + z = \text{const.}$  ... (4)

and  $x dx + y dy + z dz = 0$  gives  $x^2 + y^2 + z^2 = \text{const.}$  ... (5)

(4) and (5) together represent the curves of equal pressure and density. These curves being the curves of intersection of planes and spheres represent circles.

**Example 5.** A fluid rests in equilibrium in a field of force,  $X = y^2 + z^2 - xy - xz$ ,  $Y = z^2 + x^2 - zy - xy$ ,  $Z = x^2 + y^2 - xz - yz$ . Show that the curves of equal pressure and density are a set of circles.

**Sol.** The differential equations of curves of equal pressure and density are

$$\frac{dx}{\frac{\partial Y}{\partial z} - \frac{\partial Z}{\partial y}} = \frac{dy}{\frac{\partial Z}{\partial x} - \frac{\partial X}{\partial z}} = \frac{dz}{\frac{\partial X}{\partial y} - \frac{\partial Y}{\partial x}}$$

or

$$\frac{dx}{(2z - y) - (2y - z)} = \frac{dy}{(2x - z) - (2z - x)} = \frac{dz}{(2y - x) - (2x - y)}$$

or

$$\frac{dx}{y - z} = \frac{dy}{z - x} = \frac{dz}{x - y}$$

$$= \frac{dx + dy + dz}{0} = \frac{x dx + y dy + z dz}{0}$$

Now

$$dx + dy + dz = 0 \text{ gives } x + y + z = \text{const.}$$

and

$$x dx + y dy + z dz = 0 \text{ gives } x^2 + y^2 + z^2 = \text{const.}$$

These equations together represent a set of circles.

**Example 6.** If the components parallel to the axes of the forces acting on the element of fluid at  $(x, y, z)$  be proportional to

$$y^2 + 2\lambda yz + z^2, z^2 + 2\mu zx + x^2, x^2 + 2vxy + y^2,$$

show that if equilibrium be possible, we must have

$$2\lambda = 2\mu = 2v = 1.$$

[TMBU-2006H]

**Sol.** Here let  $X = k(y^2 + 2\lambda yz + z^2)$ ,

$$Y = k(z^2 + 2\mu zx + x^2), Z = k(x^2 + 2vxy + y^2).$$

The fluid will be at rest if

$$X \left( \frac{\partial Y}{\partial z} - \frac{\partial Z}{\partial y} \right) + Y \left( \frac{\partial Z}{\partial x} - \frac{\partial X}{\partial z} \right) + Z \left( \frac{\partial X}{\partial y} - \frac{\partial Y}{\partial x} \right) = 0 \quad \dots(1)$$

or if

$$k^2 \Sigma \{(y^2 + 2\lambda yz + z^2)(2z + 2\mu x - 2vx - 2y)\} = 0$$

or

$$\Sigma \{(y^2 + 2\lambda yz + z^2)\}(z - y + (\mu - v)x) = 0.$$

When  $2\lambda = 2\mu = 2v = 1$ , it becomes

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$$\Sigma \{(y^2 + yz + z^2)(z - y)\} = 0$$

or

$$\Sigma (y^3 - z^3) = 0$$

or

$$(y^3 - z^3) + (z^3 - x^3) + (x^3 - y^3) = 0$$

which is clearly satisfied.

Thus the condition of equilibrium is satisfied when

$$2\lambda = 2\mu = 2\nu = 1.$$

This proves the result.

**Example 14.** A given volume of liquid is at rest on a fixed plane under the action of a force, to a fixed point in the plane, varying as the distance. Find the pressure at any point of the liquid and the whole pressure on the fixed plane.

Sol. Let the fixed point be taken as origin. Consider a point  $P$  at a distance  $r$  from the origin. Force on it  $= \mu r$  towards the centre; hence pressure at this point is given by

$$dp = -\rho \mu r dr.$$

$$\text{Integrating, } p = c - \frac{1}{2} \rho \mu r^2.$$

And if  $\frac{2}{3} \pi a^3$  be the given volume, the free surface is a hemisphere of radius  $a$ . Thus we have

$$p = 0 \text{ when } r = a$$

$$\therefore c = \frac{1}{2} \rho \mu a^2$$

$$\therefore p = \frac{1}{2} \rho \mu (a^2 - r^2).$$

This gives pressure at any point  $P$  of the liquid.

To find whole pressure on the fixed plane. To find total pressure on the fixed plane (which is here a circle of radius  $a$ ) consider a point  $Q$  on the plane at a distance  $u$  from  $O$  enclosed by an element of area  $u \delta u \delta \theta$ .

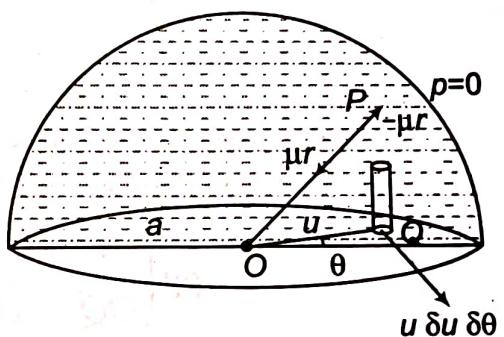


Fig. 170

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Putting  $r = u$  in (1), pressure at  $Q = \frac{1}{2} \rho \mu (a^2 - \mu^2)$ .

$$\text{Pressure on the element } u \delta u \delta \theta = \frac{1}{2} \rho \mu (a^2 - \mu^2) u \delta u \delta \theta.$$

$$\text{So total pressure} = \int_0^{2\pi} \int_0^a \frac{1}{2} \rho \mu (a^2 - u^2) u du d\theta$$

$$= \pi \rho \mu \left[ \frac{1}{2}a^2 u^2 - \frac{1}{4}u^4 \right]_0^a = \pi \rho \mu \left[ \frac{1}{2}a^4 - \frac{1}{4}a^4 \right] = \frac{1}{4}\pi \rho \mu a^4.$$

**Ex 15** \*Example 15. A mass of liquid rests upon a plane subject to a central attractive force  $\mu/r^2$ , situated at a distance  $c$  from the plane on the side opposite to that on which is the fluid; show that the pressure on the plane =  $\frac{\pi \rho \mu (a - c)^2}{a}$ .

**Sol.** The liquid rests on the plane in the form of a cap which is part of a sphere of radius  $a$ . We have to determine pressure on the plane (the circular base of the cap). Let us first consider a point  $P$  in the liquid at a distance  $r$  from the centre  $O$  of the sphere.

The only force on  $P$  is  $\frac{\mu}{r^2}$  towards  $O$ . Therefore the

pressure at  $P$  is given by  $dp = -\rho \frac{\mu}{r^2} dr$ .

Integrating  $p = c + \frac{\rho\mu}{r}$

But when  $r = a, p = 0$  at the free surface ;  $\therefore c = - \frac{\rho \mu}{a}$

$$\text{Thus } p = \rho \mu \left( \frac{1}{r} - \frac{1}{a} \right) = \rho \mu \left( \frac{1}{OP} - \frac{1}{a} \right)$$

Now consider a point  $Q$  on the plane base of the cap; then pressure at  $O$

$$= \rho \mu \left( \frac{1}{OO} - \frac{1}{q} \right).$$

Take an element  $u \delta u \delta\theta$  surrounding the point  $Q$  such that  $CQ = u$  and  $\angle COQ = \theta$ , so that  $OQ^2 = OC^2 + CQ^2 = c^2 + u^2$ ; then pressure on the small element  $u \delta u \delta\theta$  at  $Q$

$$\rho\mu \left( \frac{1}{OQ} - \frac{1}{a} \right) + u \delta u \delta \theta = \rho\mu \left\{ \frac{1}{\sqrt{(c^2 + u^2)}} - \frac{1}{a} \right\} u \delta u \delta \theta.$$

Integrating between the proper limits, so as to include the whole circular base of the cap. total pressure on the plane.

$$= \rho \mu \int_0^{\sqrt{(a^2 - c^2)}} \int_0^{2\pi} \left\{ \frac{1}{\sqrt{(c^2 + u^2)}} - \frac{1}{a} \right\} u \, du \, d\theta,$$

[radius of the base being  $\sqrt{a^2 - c^2}$ ]

$$= \rho \mu [\theta]_0^{2\pi} \left[ \sqrt{(c^2 + u^2)} - \frac{u^2}{2a} \right] \sqrt{(a^2 - c^2)}$$

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$$\begin{aligned} &= 2\pi\rho\mu \left( a - c - \frac{a^2 - c^2}{2a} \right) \\ &= 2\pi\rho\mu (a - c) \left( 1 - \frac{a + c}{2a} \right) = \frac{\pi\rho\mu}{a} (a - c)^2. \end{aligned}$$